Fujimoto-Watanabe equations and differential substitutions

This article has been downloaded from IOPscience. Please scroll down to see the full text article.
1991 J. Phys. A: Math. Gen. 24 L519
(http://iopscience.iop.org/0305-4470/24/10/004)
View the table of contents for this issue, or go to the journal homepage for more

Download details:
IP Address: 129.252.86.83
The article was downloaded on 01/06/2010 at 10:23

Please note that terms and conditions apply.

## LETTER TO THE EDITOR

# Fujimoto-Watanabe equations and differential substitutions 

S Yu Sakovich<br>Institute of Physics, BSSR Academy of Sciences, Minsk 220602, USSR

Received 25 February 1991


#### Abstract

New Fujimoto-Watanabe non-constant separant evolution equations admitting generalized symmetries are shown to be connected with the Korteweg-de Vries, SawadaKotera, Kaup and linear equations via chains of differential substitutions. A conjecture is proposed which explains why the Ibragimov transformation to constant separant equations is a universal link in those chains.


Recently, Fujimoto and Watanabe [1] classified third-order polynomial evolution equations of uniform rank with non-constant separants, which admit generalized symmetries [2] (non-trivial Lie-Bäcklund algebras [3]), and found eight equations of that kind:

$$
\begin{align*}
& u_{t}=u_{1}^{3} u_{3}+\alpha u_{1}^{4}  \tag{1}\\
& u_{t}=u_{1}^{3} u_{3}+\alpha u_{1}^{3}  \tag{2}\\
& u_{t}=u^{3} u_{3}+3 u^{2} u_{1} u_{2}+\alpha\left(u^{3} u_{2}+u^{2} u_{1}^{2}\right)+\frac{2}{9} \alpha^{2} u^{3} u_{1}  \tag{3}\\
& u_{t}=u^{3} u_{3}+3 u^{2} u_{1} u_{2}+4 \alpha u^{3} u_{1}  \tag{4}\\
& u_{t}=u^{3} u_{3}+3 u^{2} u_{1} u_{2}+3 \alpha u^{2} u_{1}  \tag{5}\\
& u_{t}=u^{3} u_{3}+\alpha u^{3} u_{1}  \tag{6}\\
& u_{t}=u^{3} u_{3}+\frac{3}{2} u^{2} u_{1} u_{2}+\alpha\left(u^{3} u_{2}+u^{2} u_{3}^{2}\right)+\frac{2}{9} \alpha^{2} u^{3} u_{1}  \tag{7}\\
& u_{t}=u^{3} u_{3}+\frac{3}{2} u^{2} u_{1} u_{2}+\alpha u^{2} u_{1} \tag{8}
\end{align*}
$$

where $u=u(x, t), u_{t}=\partial u / \partial t, u_{k}=\partial^{k} u / \partial x^{k}, \alpha$ is a parameter. Fujimoto and Watanabe found also two fifth-order equations

$$
\begin{align*}
& u_{t}=u^{5} u_{5}+5 u^{4}\left(u_{1} u_{4}+2 u_{2} u_{3}\right)  \tag{9}\\
& u_{t}=u^{5} u_{5}+5 u^{4}\left(u_{1} u_{4}+\frac{1}{2} u_{2} u_{3}\right)+\frac{15}{4} u^{3} u_{1}^{2} u_{3} \tag{10}
\end{align*}
$$

which also admit generalized symmetries and do not belong to hierarchies of equations (1)-(8). Fujimoto and Watanabe pointed out that (2) and (1) are the potential forms of (5) and (4), the last two being equivalent to the Korteweg-de Vries (Kdv) equation and the modified KdV equation, respectively, via the Ishimori transformation defined as $x=\int u \mathrm{~d} x$ in [1].

In this letter, we construct chains of differential substitutions which connect the Fujimoto-Watanabe equations (1)-(10) with the Kdv, Sawada-Kotera [4], Kaup [5] and linear equations. We consider (1)-(10), including equations transformed into the KdV equation by Fujimoto and Watanabe, in order to demonstrate and generalize certain common features of the chains of differential substitutions involved. Proofs, which can be restored easily, are omitted here.

Equations (3)-(10) have the form $u_{t}=u^{n} u_{n}+\ldots, n=3$ or 5 . As one can prove, any differential substitution (if it exists) connecting equations $u_{r}=u^{n} u_{n}+\ldots$ and $v_{r}=$ $v_{n}+\ldots(n>1)$ has the form $x=b, u=D_{y} b$, where $b\left(y, v, v_{1}, \ldots, v_{m}\right)$ is a certain function, $v=v(y, t), v_{k}=\partial^{k} v / \partial y^{k}, D_{y}$ is the total $y$-derivative. Such a substitution with $b=v$ was introduced by Ibragimov [3,6] as a transformation between the Harry Dym equation and modified via the Schwarzian derivative KdV equation (respectively (6) above and (16) below at $\alpha=0$ ). The lbragimov transformation $x=v, u=v_{1}$ can be proven to map only the equations

$$
\begin{equation*}
v_{t}=\xi y v_{1}+v_{1} a\left(v, v_{1}, v_{1}^{-1} v_{2}, \ldots,\left(v_{1}^{-1} D_{y}\right)^{n-1} v_{1}\right) \tag{11}
\end{equation*}
$$

into the corresponding ones

$$
\begin{equation*}
u_{t}=\xi u+u^{2} D_{x} a\left(x, u, u_{1}, \ldots, u_{n-1}\right) \tag{12}
\end{equation*}
$$

with arbitrary parameter $\xi$, function $a$ and order $n$ (the proof is based on the technique of $[7,8]$ ). Equations (3)-(10) actually have the form $u_{t}=u^{2} D_{x}(\ldots)$. This is a surprising phenomenon but not a chance one. Indeed, any equation $u_{t}=f\left(x, u, u_{1}, \ldots, u_{n}\right)$ admitting a generalized symmetry of order higher than $n$ has the conserved density $\left(\partial f / \partial u_{n}\right)^{-1 / n}$ [3]. If $f=u^{n} u_{n}+\ldots$, then $\left(u^{-1}\right)_{t}+D_{x}(\ldots)=0$ takes place, and $f=$ $u^{2} D_{x}(\ldots)$ turns out to be necessary. Therefore (3)-(10) belong to class (12), and corresponding equations of class (11) are as follows:

$$
\begin{align*}
& v_{\mathrm{t}}=v_{3}+\alpha v_{1} v_{2}+\frac{1}{9} \alpha^{2} v_{1}^{3}  \tag{13}\\
& v_{\mathrm{t}}=v_{3}+2 \alpha v_{1}^{3}  \tag{14}\\
& v_{\mathrm{t}}=v_{3}+3 \alpha v_{1}^{2}  \tag{15}\\
& v_{\mathrm{t}}=v_{3}-\frac{3}{2} v_{1}^{-1} v_{2}^{2}+\frac{1}{2} \alpha v_{1}^{3}  \tag{16}\\
& v_{\mathrm{t}}=v_{3}-\frac{3}{4} v_{1}^{-1} v_{2}^{2}+\alpha v_{1} v_{2}+\frac{1}{9} \alpha^{2} v_{1}^{3}  \tag{17}\\
& v_{\mathrm{t}}=v_{3}-\frac{3}{4} v_{1}^{-1} v_{2}^{2}+\alpha v_{1}^{2}  \tag{18}\\
& v_{t}=v_{5}-5 v_{1}^{-1} v_{2} v_{4}+5 v_{1}^{-2} v_{2}^{2} v_{3}  \tag{19}\\
& v_{t}=v_{5}-5 v_{1}^{-1} v_{2} v_{4}-\frac{15}{4} v_{1}^{-1} v_{3}^{2}+\frac{65}{4} v_{1}^{-2} v_{2}^{2} v_{3}-\frac{135}{16} v_{1}^{-3} v_{2}^{4} \tag{20}
\end{align*}
$$

The correspondence under the Ibragimov transformation $x=v, u=v_{1}$ is $(k+10) \rightarrow(k)$, $k=3,4, \ldots, 10$. In the right-hand sides of (13)-(20) we omit the trivial term $\beta v_{1}$ caused by the invariance of (12) under $a \rightarrow a+\beta, \beta$ is any constant.

Now, constant separant equations (13)-(20) can be reduced to more familiar ones. Third-order equations (13)-(18) fall under Habirov's classification [9] which provides the following: (13) and (17) are mapped into the linear equation $w_{1}=w_{3}$ via $w=\exp \left(\frac{1}{3} \alpha v\right)$ and $w=\left(\frac{2}{3} \alpha v_{1}\right)^{1 / 2} \exp \left(\frac{1}{3} \alpha v\right)$ respectively; (14), (15), (16) and (18) are mapped onto the Kdv equation $w_{1}=w_{3}+3 w w_{1}$ via $w=2(-\alpha)^{1 / 2} v_{2}+2 \alpha v_{1}^{2}, w=2 \alpha v_{1}$, $w=v_{1}^{-1} v_{3}-\frac{3}{2} v_{1}^{-2} v_{2}^{2}+\frac{1}{2} \alpha v_{1}^{2}$ and $w=\left(-\frac{1}{3} \alpha\right)^{1 / 2} v_{1}^{-1 / 2} v_{2}+\frac{2}{3} \alpha v_{1}$ respectively; $w=w(y, t)$, $w_{k}=\partial^{k} w / \partial y^{k}$. Fifth-order equations (19) and (20) are mapped onto the Kaup equation [5] $w_{t}=w_{5}+10 w w_{3}+25 w_{1} w_{2}+20 w^{2} w_{1}$ and Sawada-Kotera equation [4] $w_{t}=$ $w_{5}+\frac{5}{3} w w_{3}+\frac{5}{2} w_{1} w_{2}+\frac{5}{4} w^{2} w_{1}$, respectively, via $w=v_{1}^{-1} v_{3}-\frac{3}{2} v_{1}^{-2} v_{2}^{2}$ [7] which is the Schwarzian derivative. Equation (19) is also mapped onto the Sawada-Kotera equation via the differential substitution $w=-2 v_{1}^{-1} v_{3}$ associated with the inverse scattering transform technique [10] at zero spectral parameter.

The remaining equations (1) and (2) admit the map $z=u_{1}$ into (4) and (5) on $z$ respectively [1]. Note that the form $z=\left(\partial f / \partial u_{n}\right)^{1 / n}$ is necessary for any differential substitution which maps an equation $u_{t}=f\left(x, u, u_{1}, \ldots, u_{n}\right)$ onto $z_{t}=z^{n} z_{n}+\ldots$ Does this substitution exist for every non-constant separant evolution equation admitting a generalized symmetry? Is the Ibragimov transformation applicable to the resulting equation $z_{t}=z^{n} z_{n}+\ldots$ ? Do the involved transformations preserve generalized symmetries? We believe the following conjecture answers these questions correctly.

Conjecture. If two non-constant separant evolution equations $u_{t}=f\left(x, u, u_{1}, \ldots, u_{n}\right)$ and $u_{s}=g\left(x, u, u_{1}, \ldots, u_{m}\right) \quad(m>n>1)$ are compatible, i.e. $\Sigma_{k}\left(\partial f / \partial u_{k}\right) D_{x}^{k} g \equiv$ $\Sigma_{k}\left(\partial g / \partial u_{k}\right) D_{x}^{k} f$, then: (i) these equations admit the differential substitution $z=$ $\left(\partial f / \partial u_{n}\right)^{1 / n} \equiv\left(\partial g / \partial u_{m}\right)^{1 / m}$ (scaling of $s$ is implied); (ii) the resulting equations $z_{y}=$ $z^{n} z_{n}+\ldots$ and $z_{s}=z^{m} z_{m}+\ldots$ are compatible and belong to class (12); (iii) the corresponding equations $v_{t}=v_{n}+\ldots$ and $v_{s}=v_{m}+\ldots$ of class (11) are compatible too.

At present, we have a fragmentary proof only, and the conjecture should be used as certain tactical means. In addition to the Fujimoto-Waranabe equations, let us quote one more example illustrating the conjecture in action. The equation

$$
\begin{equation*}
u_{t}=D_{x}\left(u^{-3} u_{2}-3 u^{-4} u_{1}^{2}\right)+1 \tag{21}
\end{equation*}
$$

admits an infinite algebra of generalized symmetries which is destroyed by the transformation $y=u, w=u_{1}[11]$. Our conjecture indicates transformations which preserve generalized symmetries and connect (21) with the KdV equation. Indeed, the separant of (21) is $u^{-3}$, the substitution $z=u^{-1}$ gives $z_{t}=z^{2} D_{x}\left(z z_{2}+z_{1}^{2}-x\right)$ which belong to class (12), and the corresponding equation of class (11) is $v_{t}=v_{3}-v v_{1}$.

Regardless of the validity of our conjecture, differential substitutions need to be used in any classification of nonlinear partial differential equations. Anyhow, every new remarkable equation must be suspected of being an appropriately spoilt well known one.

## References

[1] Fujimoto A and Watanabe Y 1989 Phys. Lett. 136A 294-9
[2] Olver P J 1986 Applications of Lie Groups to Differential Equations (New York: Springer)
[3] Ibragimov N H 1985 Transformation Groups Applied to Mathematical Physics (Dordrecht: Reidel)
[4] Sawada K and Kotera T 1974 Prog. Theor. Phys. 51 1355-67
[5] Kaup D J 1980 Stud. Appl. Math. 62 189-216
[6] Ibragimov N H 1981 C.R. Acad. Sci. Paris 293 657-60
[7] Sakovich S Yu 1988 Phys. Lett. 132A 9-12
[8] Sakovich S Yu 1990 Phys. Lett. 146A 32-4
[9] Habirov S V 1985 Zh. Vych. Mat. Mat. Fiz. 25 935-41
[10] Dodd R K and Gibbon J D 1977 Proc. R. Soc. A 358 287-96
[11] Sokolov V V 1988 Usp. Mat. Nauk 43 133-63

